電子回路論第8回 Electric Circuits for Physicists #8

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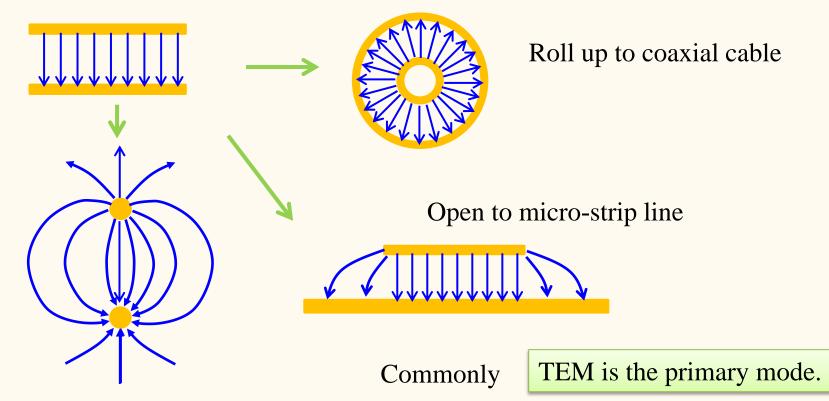
Shingo Katsumoto

Outline

Various transmission lines Termination and connection Smith chart S-matrix (S-parameters)

Transmission lines with TEM mode

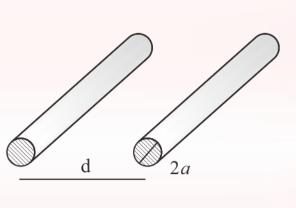
Transmission lines with two conductors are "families". Electromagnetic field confinement with parallel-plate capacitor

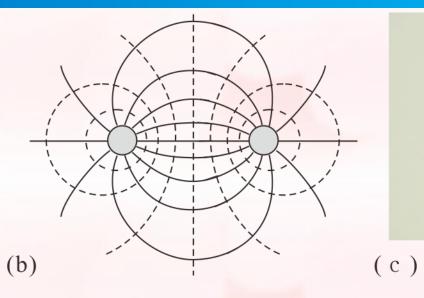


Shrink to dipole (Lecher line)

Lecher line

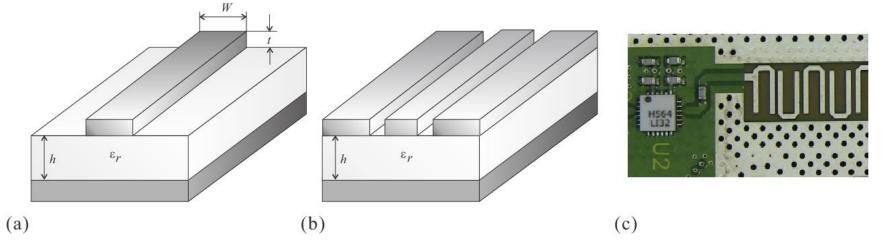
(a)





$$\phi_1 = -\phi_2 = \frac{J\sqrt{\mu}}{2\pi}\log\frac{d}{a}$$
 $Z_0 = \sqrt{\frac{\mu}{\epsilon}}\frac{1}{\pi}\log\frac{d}{a}$

Micro strip line



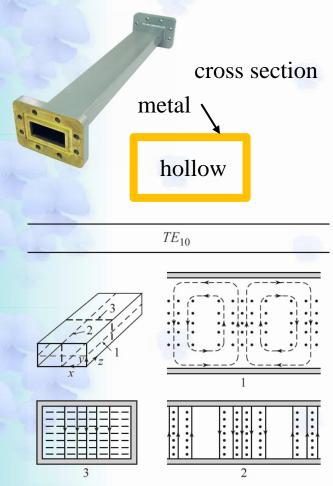
Wide (W/h>3.3) strip

$$Z(W,h,\epsilon_r) = \frac{Z_{F0}}{2\sqrt{\epsilon_r}} \left\{ \frac{W}{2h} + \frac{1}{\pi} \log 4 + \frac{\epsilon_r + 1}{2\pi\epsilon_r} \log \left[\frac{\pi e}{2} \left(\frac{W}{2h} + 0.94 \right) \right] \frac{\epsilon_r - 1}{2\pi\epsilon_r^2} \log \frac{e\pi^2}{16} \right\}^{-1}$$

Narrow (W/h<3.3) strip

$$Z(W,h,\epsilon_r) = \frac{Z_{F0}}{\pi\sqrt{2(\epsilon_r+1)}} \left\{ \log\left[\frac{4h}{W} + \sqrt{\left(\frac{4h}{W}\right)^2 + 2}\right] - \frac{1}{2}\frac{\epsilon_r - 1}{\epsilon_r + 1} \left(\log\frac{\pi}{2} + \frac{1}{\epsilon_r}\log\frac{4}{\pi}\right) \right\}$$

Waveguide



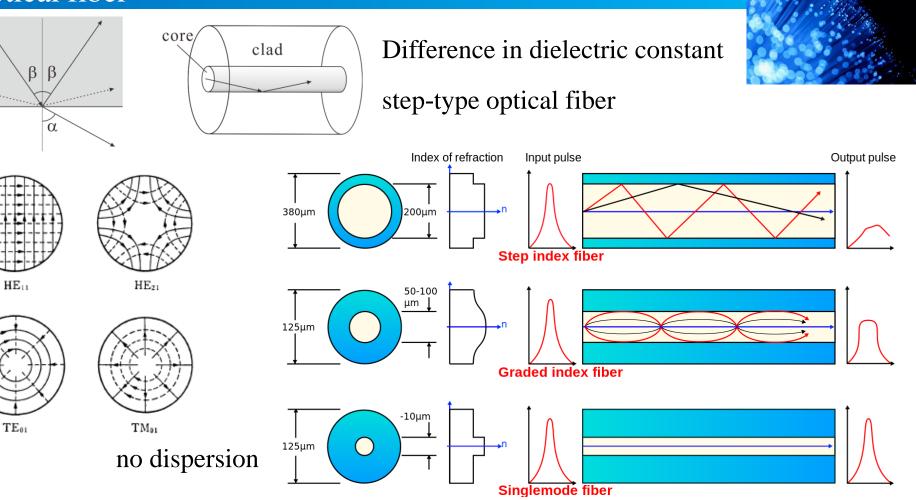
Electromagnetic field is confined into a simplyconnected space. TEM mode cannot exist. Maxwell equations give $\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right] E_z = -(\omega^2 \epsilon \mu + \gamma^2) E_z,$ $\left[\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right]H_z = -(\omega^2\epsilon\mu + \gamma^2)H_z.$ Helmholtz equation $E_{z} = 0$: TE mode,

 $H_z = 0$: TM mode

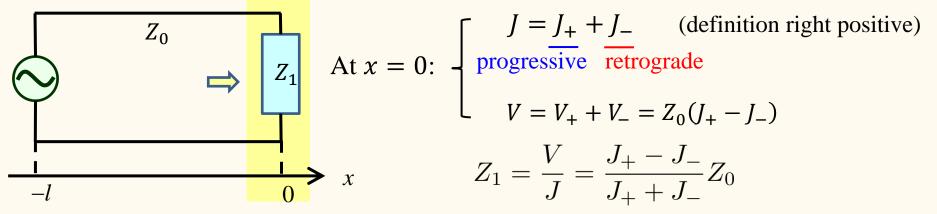
Optical fiber

 n_2

 n_{1}



Connection and termination



Termination of a transmission line with length *l* and characteristic impedance Z_0 at x = 0 with a resistor Z_1 .

Comment: Sign of Ohm's law in transmission lines reflects direction of waves (depends on the definitions).

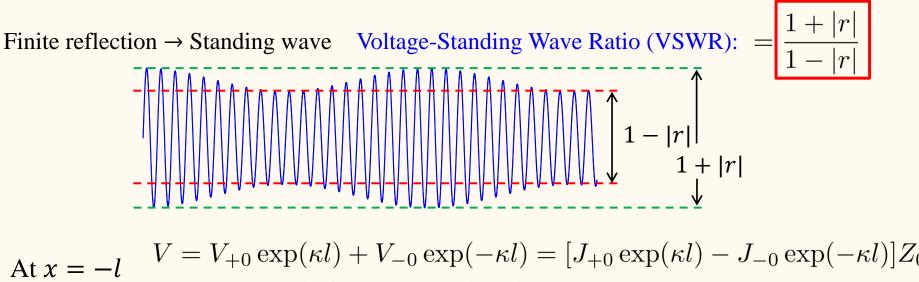
Reflection coefficient is $r = \frac{V_{-}}{V_{+}} = -\frac{J_{-}}{J_{+}} = \frac{Z_{1} - Z_{0}}{Z_{1} + Z_{0}}$

 $Z_1 = Z_0$: no reflection, *i.e.*, impedance matching

 $Z_1 = +\infty$ (open circuit end) : r = 1, *i.e.*, free end

 $Z_1 = 0$ (short circuit end) : r = -1, *i.e.*, fixed end

Connection and termination



At
$$x = -l$$
 $V = V_{+0} \exp(\kappa l) + V_{-0} \exp(-\kappa l) = [J_{+0} \exp(\kappa l) - J_{-0} \exp(-\kappa l)]Z_0$
 $J = J_{+0} \exp(\kappa l) + J_{-0} \exp(-\kappa l)$

Then at x = -l (at power source), the right hand side can be represented by

$$Z_{l} = \frac{V}{J} = \frac{J_{+0}e^{\kappa l} - J_{-0}e^{-\kappa l}}{J_{+0}e^{\kappa l} + J_{-0}e^{-\kappa l}}Z_{0}$$

Reflection coefficient:
$$r_l = \frac{V_-}{V_+} = \frac{V_{0-} \exp(-\kappa l)}{V_{0+} \exp(\kappa l)} = r \exp(-2\kappa l)$$

Connection and termination



 Z_0

Transmission line connection. Characteristic impedance Z_0, Z_0'

At the connection point, only the local relation between V and J affects the reflection coefficient.

The local impedance from the left hand side is Z_0' .

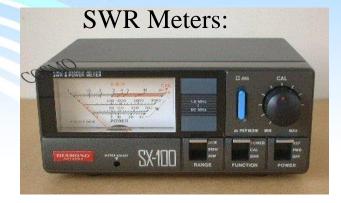
$$r = \frac{Z_0' - Z_0}{Z_0' + Z_0}$$

antenna

$$Z_0^{\rm v} = 120\pi \approx 377 \ \Omega$$

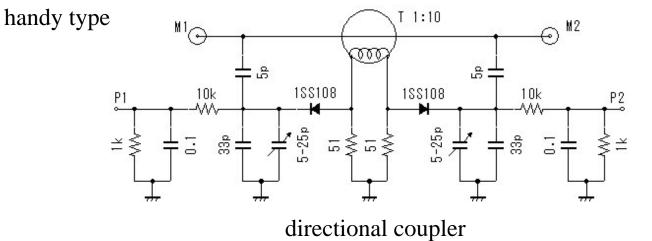
 Z'_0 : input impedance of antenna

SWR measurement



desktop types cross-meter





5.2.3 Smith chart, immittance chart

End impedance Z_1 : Normalized end impedance $Z_n \equiv Z_1/Z_0$

$$Z_{n} = x + iy, \quad r = u + iw \quad (x, y, u, w \in \mathbb{R})$$

$$u + iw = r = \frac{Z_{n} - 1}{Z_{n} + 1} = \frac{(x - 1) + iy}{(x + 1) + iy} \qquad \text{real:} \qquad x - 1 = (x + 1)u - yw$$

$$\text{imaginary:} \qquad y = yu + w(x + 1)$$

$$\left(u - \frac{x}{x+1}\right)^2 + w^2 = \frac{1}{(x+1)^2}$$

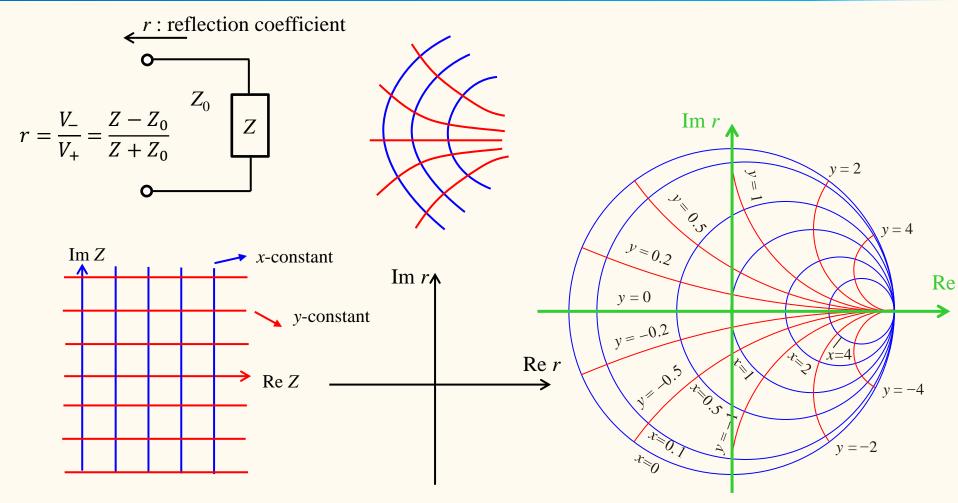
constant reactance circle

x: erase *y*: constant

$$(u-1)^2 + \left(w - \frac{1}{y}\right)^2 = \frac{1}{y^2}$$

constant reactance circle

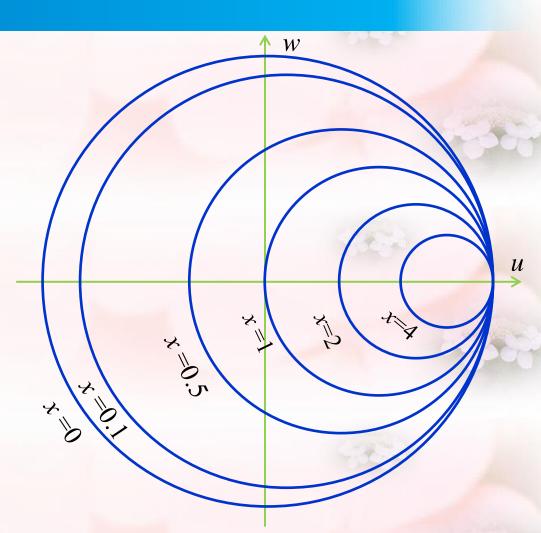
5.2.3 Smith chart, immittance chart



Circles with constant *x*

$$r = u + iw$$
$$|r| \le 1$$
$$\left(u - \frac{x}{x+1}\right)^2 + w^2 = \frac{1}{(x+1)^2}$$

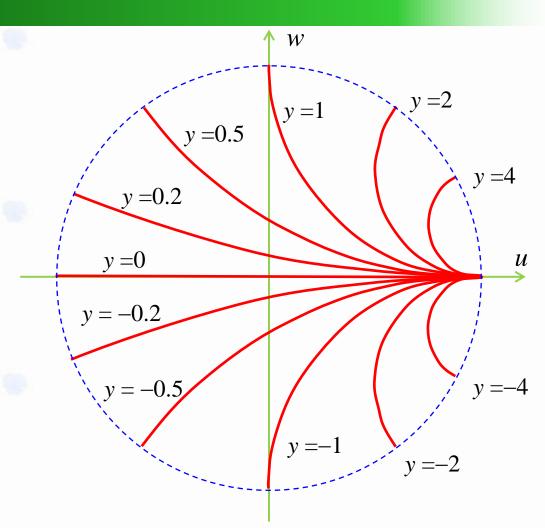
Common point: (u, w) = (1, 0)



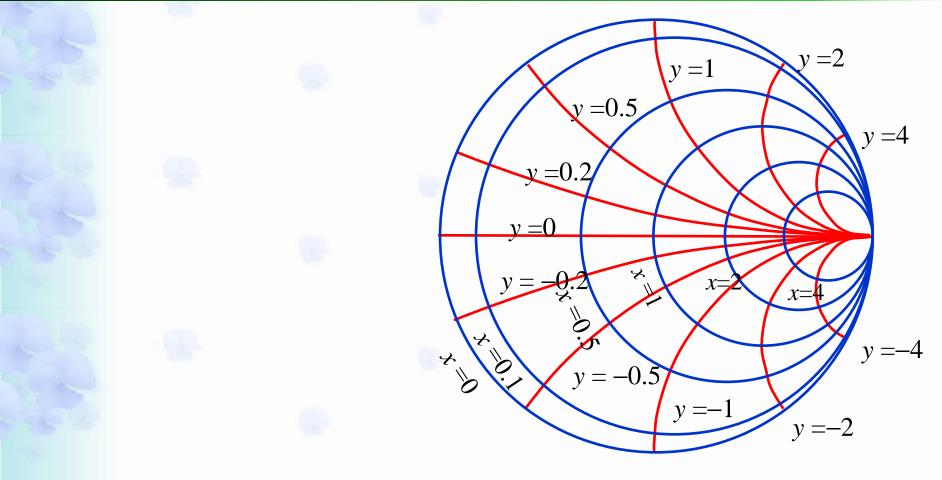
Circles with constant *y*

$$r = u + iw$$
$$|r| \le 1$$
$$(u-1)^2 + \left(w - \frac{1}{y}\right)^2 = \frac{1}{y^2}$$

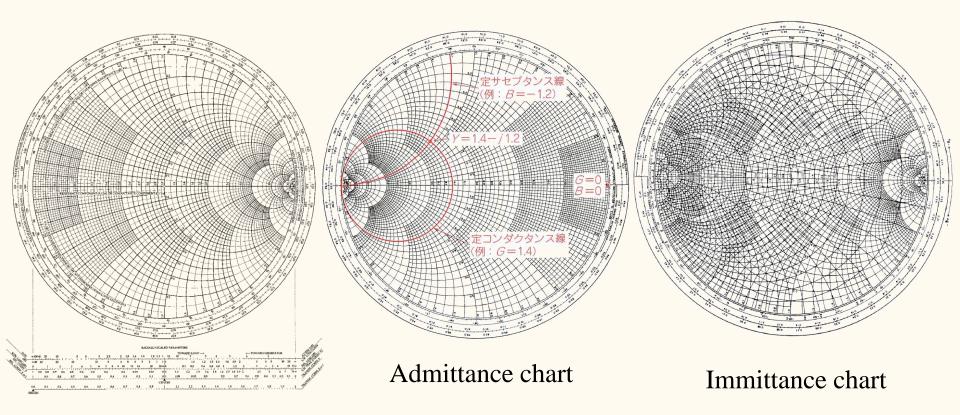
Common point: (u, w) = (1, 0)



Smith chart

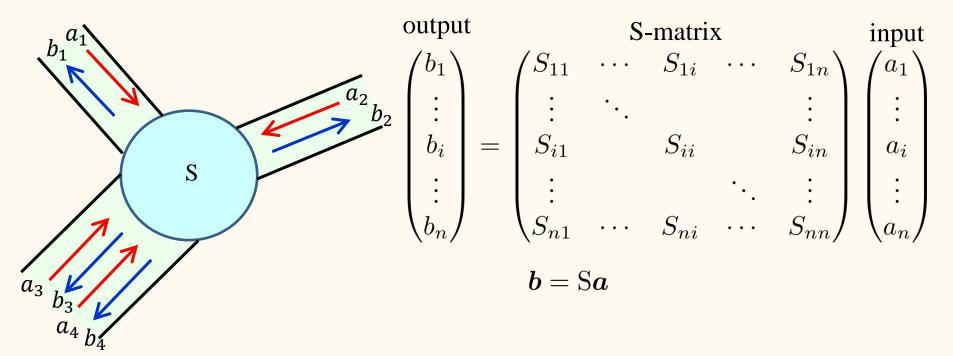


Smith chart, immittance chart



5.3 Scattering (S) matrix (S parameters)

How to treat multipoint (crossing point) in transmission lines systematically? Transmission lines: wave propagating modes \rightarrow Channels Take $|a_i|^2$, $|b_i|^2$ to be powers (energy flow).



S-matrix symmetries

(In case no dissipation, no amplification)

Propagation with no dissipation

$$\begin{cases} a_n = \frac{V_{n+}}{\sqrt{Z_{0n}}} = J_{n+}\sqrt{Z_{0n}}, \\ b_n = \frac{V_{n-}}{\sqrt{Z_{0n}}} = J_{n-}\sqrt{Z_{0n}}, \end{cases}$$

Reciprocity
$$S_{ij} = S_{ji}$$

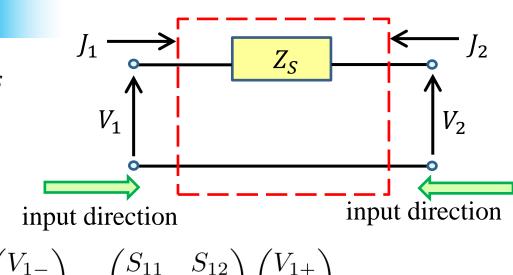
Unitarity $\sum_{j} S_{ji} S_{jk}^* = \delta_{ik}$
, Incident power wave
Reflected (transmitted)
power wave

$$|a_n|^2 = \frac{|V_{n+}|^2}{Z_{0n}} = |J_{n+}|^2 Z_{0n}$$

S matrix (S parameters)

Simplest example: series impedance Z_S

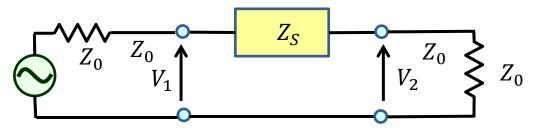
Take voltage as the "flow" quantity. (assume common characteristic impedance)



$$\begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \begin{pmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} \qquad \begin{pmatrix} V_{1-} \\ V_{2-} \end{pmatrix} = \begin{pmatrix} S_{11} & S_{12} \\ S_{21} & S_{22} \end{pmatrix} \begin{pmatrix} V_{1+} \\ V_{2+} \end{pmatrix}$$

Measurement (calculation) of the S-matrix in the case the above is placed in transmission line with characteristic impedance Z_0 .

Consider the case that the ends of the line is terminated with Z_0 and the source is connected at the left end.



Simple example of S-matrix

Terminate 2 with $Z_0 \rightarrow a_2 = 0$ (no reflection, no wave source at the right end)

$$S_{11} = \frac{V_{1-}}{V_{1+}} = \frac{V_1 - Z_0 J_1}{V_1 + Z_0 J_1} = \frac{Z_1 - Z_0}{Z_1 + Z_0} = \frac{(Z_S + Z_0) - Z_0}{(Z_S + Z_0) + Z_0} = \frac{Z_S}{Z_S + 2Z_0}$$
$$S_{21} = \frac{V_{2-}}{V_{1+}} = \frac{V_2 - Z_0 J_2}{V_1 + Z_0 J_1} = \frac{Z_0 J_1 + Z_0 J_1}{(Z_S + Z_0) J_1 + Z_0 J_1} = \frac{2Z_0}{Z_S + 2Z_0} (J_2 = -J_1)$$

From symmetry of S-matrix: $S_{22} = S_{11} = \frac{Z_S}{Z_S + 2Z_0}$, $S_{12} = S_{21} = \frac{2Z_0}{Z_S + 2Z_0}$

When the impedance matrix $Z = \{Z_{ij}\}$ is given, the S-matrix is obtained as

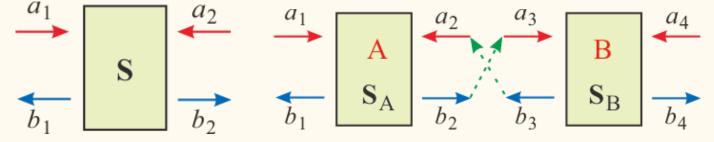
$$S = \frac{1}{\det Z} \begin{pmatrix} (Z_{11} - Z_0)(Z_{22} + Z_0) - Z_{12}Z_{21} & 2Z_0Z_{12} \\ 2Z_0Z_{21} & (Z_{11} - Z_0)(Z_{22} + Z_0) - Z_{12}Z_{21} \end{pmatrix}$$

Cascade connection of S-matrices

$$\begin{pmatrix} b_1 \\ b_2 \end{pmatrix} = \mathbf{S} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix} = \begin{pmatrix} r_{\mathrm{L}} & t_{\mathrm{R}} \\ t_{\mathrm{L}} & r_{\mathrm{R}} \end{pmatrix} \begin{pmatrix} a_1 \\ a_2 \end{pmatrix}$$

 $r_{L,R}$, $t_{L,R}$: complex reflection, transmission coefficients satisfying

$$T_{\rm L,R} = |t_{\rm L,r}|^2 = 1 - R_{\rm L,R} = 1 - |r_{\rm L,R}|^2$$

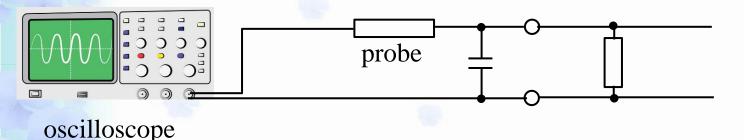


$$\mathbf{S}_{AB} = \begin{pmatrix} r_{L}^{AB} & t_{R}^{AB} \\ t_{L}^{AB} & r_{R}^{AB} \end{pmatrix} = \begin{pmatrix} r_{L}^{A} + t_{R}^{A} r_{L}^{B} \left(I - r_{R}^{A} r_{L}^{B} \right)^{-1} t_{L}^{A} & t_{R}^{A} \left(I - r_{L}^{B} r_{R}^{A} \right)^{-1} t_{R}^{B} \\ t_{L}^{B} \left(I - r_{R}^{A} r_{L}^{B} \right)^{-1} t_{L}^{A} & r_{R}^{B} + t_{L}^{B} \left(I - r_{R}^{A} r_{L}^{B} \right)^{-1} r_{R}^{A} t_{R}^{B} \end{pmatrix}$$

$$(I - r_R^A r_L^B)^{-1} = I + r_R^A r_L^B + (r_R^A r_L^B)^2 + \cdots$$

Why we need S-matrix (S-parameters) for high frequency lines?

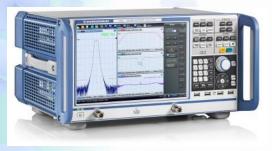
Difficulty in measurement of voltage at high frequencies.



1. Cannot ignore distributed capacitance

2. Probe line should also be treated as a transmission line

Vector network analyzer



On the other hand, S-parameters can be obtained from power measurements.

S-parameters give enough information for designing high frequency circuits.

Application of S-matrix to quantum transport

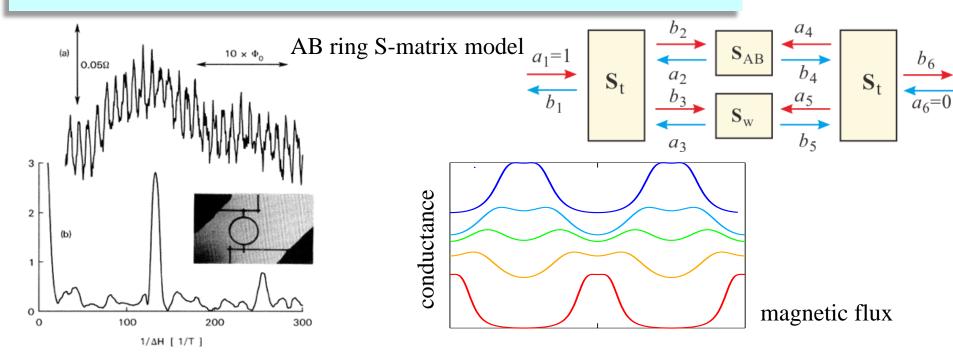
Electron (quantum mechanical) waves also have propagating modes in solids. \rightarrow Conduction channel

Landauer equation: the conductance of a single perfect quantum channel is

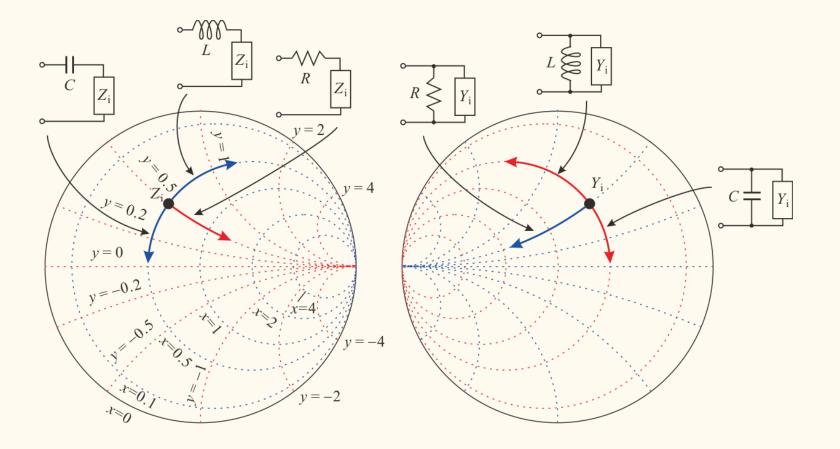


Rolf Landauer

 $\frac{e^2}{h}$



Practical impedance matching with Smith chart

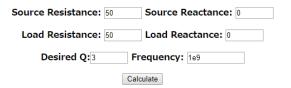


Impedance matching designer

http://home.sandiego.edu/~ekim/e194rfs01/jwmatcher/matcher2.html

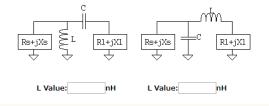
Impedance Matching Network Designer

(Now with 16 networks!)

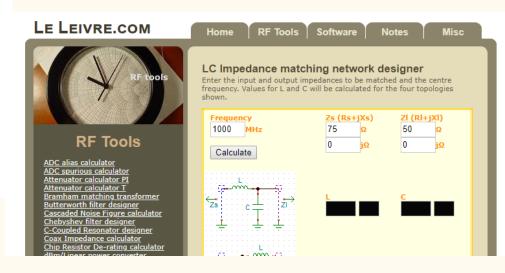


Please send comments and questions to John Wetherell at wetherel@eecs.berkeley.edu

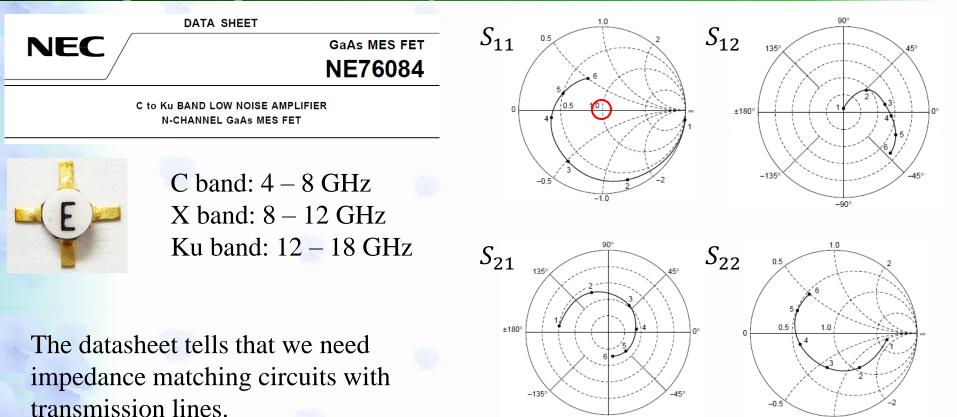
HIGHPASS Hi-Low MATCHING NETWORK LOWPASS Hi-Low MATCHING NETWORK



http://leleivre.com/rf_lcmatch.html



An example FET power matching

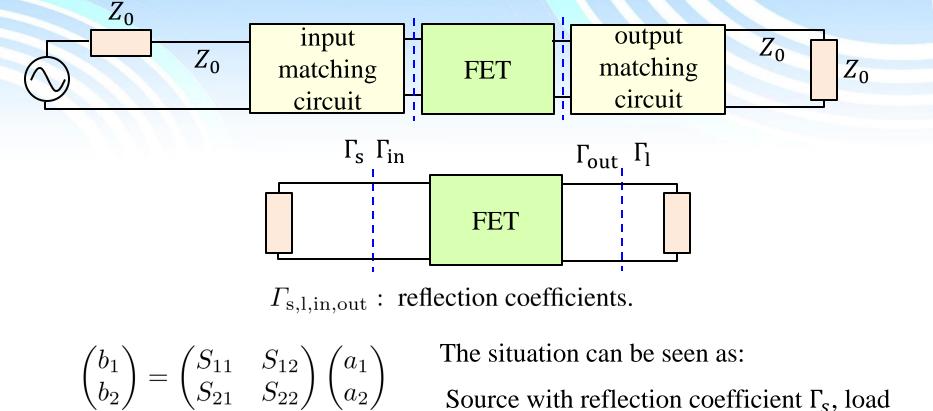


-90°

-1.0

1.25

S parameter representation and matching circuits



FET S-matrix

with reflection coefficient Γ_l are connected.

Power matching with S-parameters

That is $\Gamma_s \xrightarrow{a_1} S \xrightarrow{a_2} \Gamma_l \xrightarrow{a_1} S_l \xrightarrow{a_2} \Gamma_l \xrightarrow{a_1} S_l \xrightarrow{a_2} \Gamma_l$ Then the total coefficients are $\Gamma_{in} = S_{11} + \frac{S_{12}S_{21}\Gamma_l}{1 - S_{22}\Gamma_l}, \quad \Gamma_{out} = S_{22} + \frac{S_{12}S_{21}\Gamma_s}{1 - S_{22}\Gamma_s}$

The power matching conditions are $\Gamma_{l} = \Gamma_{out}^{*}, \quad \Gamma_{s} = \Gamma_{in}^{*}$

The quadratic equation gives $\Gamma_{\rm s} = \frac{B_1 \pm \sqrt{B_1^2 - 4|M|^2}}{2M}, \quad \Gamma_l = \frac{B_2 \pm \sqrt{B_2^2 - 4|N|^2}}{2N}$

 $B_1 = 1 + |S_{11}|^2 - |S_{22}|^2 - |\det S|^2, \quad B_2 = 1 - |S_{11}|^2 + |S_{22}|^2 - |\det S|^2,$ $N = S_{22} - S_{11}^* \det S, \quad M = S_{11} - S_{22}^* \det S$

Then the problem is reduced to tune the passive circuits to Γ_s and Γ_l