

# Physics of Semiconductors (Appendix)

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## Appendix K: Confinement with a triangular potential

To solve a problem in final report, the following knowledge on quantum confinement with a triangular potential is required.

Let us approximate a confinement potential for two-dimensional electrons with a downward triangular potential. This is a simple but a good approximation for confinement at MOS interfaces. The problem can be viewed as a slope with a limit wall, which is a well known problem in elementary quantum mechanics. We take it as a one-dimensional potential problem along  $x$ -axis and write down the Schrödinger equation as

$$\left(-\frac{\hbar^2}{2m}\frac{\partial^2}{\partial x^2} + V(x)\right)\psi = E\psi, \quad V(x) = \begin{cases} ax & (x > 0, \quad a > 0) \\ \infty & (x \leq 0) \end{cases}. \quad (8.1)$$

The equation can be rewritten with a variable transformation

$$s = \left(\frac{2ma}{\hbar^2}\right)^{1/3} \left(x - \frac{E}{a}\right). \quad (8.2)$$

as

$$\frac{d^2\psi}{ds^2} = s\psi. \quad (8.3)$$

This differential equation is mathematically called Airy or Stokes-type differential equation. The solutions are called **Airy functions**, which are classified with the asymptotic behavior for  $s \rightarrow \infty$ , into Ai ( $\psi \rightarrow 0$ ) and Bi ( $\psi \rightarrow \infty$ ). They are plotted in Fig.8.1(b). As the basis of quantum mechanical wavefunctions we should adopt Ai for vanishment at the infinite point.

Asymptotic forms of Ai for  $s \rightarrow \pm\infty$  are given as

$$\text{Ai}(s) \sim \frac{1}{2\sqrt{\pi}s^{1/4}} \exp\left(-\frac{2}{3}s^{3/2}\right) \quad (s \rightarrow \infty) \quad (8.4)$$

$$\sim \frac{1}{\sqrt{\pi}|s|^{1/4}} \cos\left(\frac{2}{3}|s|^{3/2} - \frac{\pi}{4}\right) \quad (s \rightarrow -\infty). \quad (8.5)$$

In  $x < 0$ ,  $V = \infty$  and the boundary condition at  $x = 0$  is  $\psi(+0) = 0$ . From the constraint, a zero of  $\psi(x)$  should be placed at  $x = 0$ . We write zeros of Ai from the smallest with ascending order  $s_1, s_2, \dots, s_n, \dots$ . Then from eq.(8.2), the energy eigenvalue  $E_n$  correspondent to  $n$  is obtained as

$$E_n = -\left(\frac{\hbar^2 a^2}{2m}\right)^{1/3} s_n, \quad (8.6)$$

From the asymptotic form eq.(8.5), approximate values for  $s_n$  for large  $n$  are given as

$$s_n \sim -\left(\frac{3\pi(4n-1)}{8}\right)^{2/3}. \quad (8.7)$$

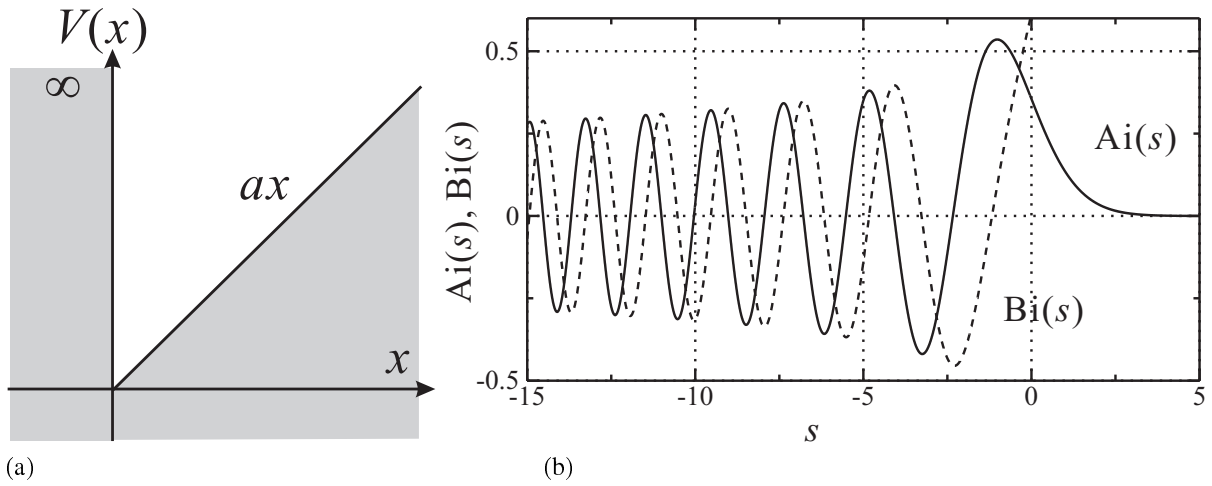


Figure 8.1: (a) Schematic illustration of a triangular potential. (b) Airy functions.

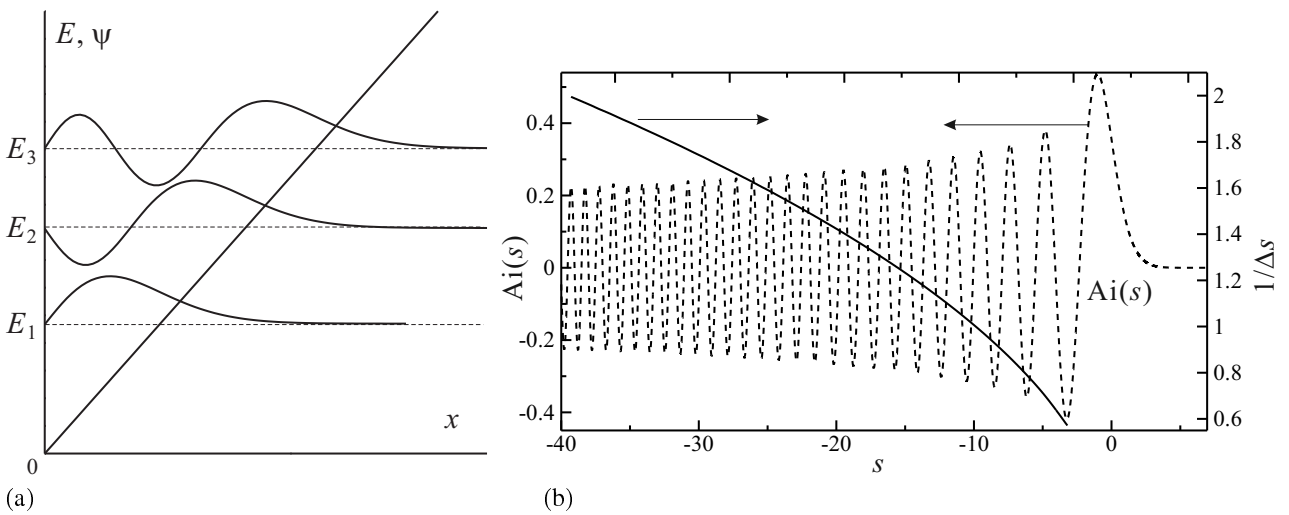


Figure 8.2: (a) Energy eigenvalues and eigenstates in a triangular confinement potential. Drawn from the lowest energy for  $n = 1, 2, 3$ . (b) Inverse of distance between neighboring zero  $\Delta s$  of Airy function  $Ai(s)$  as a function of  $s$ .  $Ai(s)$  is also plotted as a broken curve.